# Models of solar irradiance variability and the instrumental temperature record

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Abstract. The effects of decade-to-century (Dec-Cen) variations in total solar irradiance (TSI) on global mean surface temperature  $T_S$  during the pre-Pinatubo instrumental era (1854-1991) are studied by using two different proxies for TSI and a simplified version of the IPCC climate model. TSI anomalies based on solar-cycle length (CL) and solar-cycle decay rate (CD) proxies can account for most of the warming observed up to 1976, but anthropogenic forcing is needed to explain the subsequent sharp increase in  $T_S$ . The time series of CL-solar and anthropogenic radiative forcing resemble each other, making it difficult to separate their effects in the instrumental  $T_S$  record. Results using the CD-based irradiance values, however, allow tighter constraints to be placed on both TSI variability and terrestrial climate sensitivity, and underscore the inability of solar forcing alone to explain the recent global warming.

#### Introduction

Variability of the total solar irradiance (TSI) is a potentially important contributor to changes in global mean temperatures on time scales longer than a few years. Striking correlations between the instrumental  $T_S$  record, extending back nearly a century and a half, and observable solar features, such as the amplitude and length of the sunspot cycle, have suggested that solar variations may indeed have a strong impact on decade-to-century (Dec-Cen) changes in global mean temperature  $T_S$  [Reid, 1997, and references therein]. In the absence of a convincing physical link between these observed solar features and TSI, however, the role of solar variability in the terrestrial climate record is difficult to quantify. We examine here the implications of some simple physical

assumptions regarding the origin of Dec-Cen variability in TSI for the way it might affect the instrumental record of  $T_S$ .

## Models of Solar Variability

TSI variations on the order of 0.1% have been detected within a solar cycle by satellite-borne radiometers [e.g., Willson and Hudson, 1991] and successfully modeled in terms of observable photospheric features [e.g., Pap et al., 1994; Lean et al., 1998]. While their effects may be detectable in records of land surface and ocean temperatures [e.g., Stevens and North, 1996; White et al., 1997; Lawrence and Ruzmaikin, 1998], these sub-decadal fluctuations are too small to account for a significant fraction of the 0.6 °C increase in T<sub>S</sub> recorded over the last century [IPCC, 1996]. Cycle-to-cycle variability in TSI is most plausibly linked to longer-period variations in the intensity of convective heat transport from the solar interior to the photosphere [e.g., Baliunas and Jastrow, 1993; Hoyt and Schatten 1993, 1997], which may be detectable through their concomitant effects on observable solar features. In particular, Hoyt and Schatten argue that more intense convection leads to a more rapid decay of individual sunspots and a shorter solar cycle, and thus can account for the apparent (inverse) correlation between cycle length and irradiance in the sun as well as in sunlike stars [e.g. Baliunas and Soon, 1995].

We investigate here the implications of this assumption for solar effects on the instrumental temperature record. To do so, we use simplified models for cycle-to-cycle TSI variations and the response of the terrestrial mean temperature  $T_S$  to net radiative forcing. Their simplicity allows us to thoroughly test the sensitivity of results to changes in model assumptions and parameter values, and avoids the complications posed by the regionally heterogeneous nature of anthropogenic forcing and the "fingerprint" of the climate system's response [cf. Schneider, 1994].

Variations F(t) in the convective transport of heat to the photosphere from its mean are modeled as proportional to the decay rate of the solar cycle, defined as the reciprocal of the time interval between the epochs of a solar cycle maximum  $t_{\rm M}$  and the subsequent minimum  $t_{\rm m}$ :

$$F_1(t_{\rm m}) = k_1 / (t_{\rm m} - t_{\rm M});$$
 (1)

 $k_1$  is an unknown proportionality constant. The convective anomaly  $F_1$  for a cycle is defined as occurring at the time  $t_{\rm m}$  of solar minimum, when the photospheric features associated with TSI variations within a cycle are largely absent [e.g., Lean et al., 1998].

For comparison with previous studies, we also consider a formulation in which the convective heat flux anomaly is modeled as inversely proportional to the cycle length, defined as the interval between successive minima:

$$F_2(t_{m+1/2}) = k_2 / (t_{m+1} - t_m); (2)$$

here  $k_2$  is an unknown proportionality constant, and the convective anomaly  $F_2$  is defined at the time  $t_{m+1/2}$  which is midway between the epochs of successive minima  $t_m$  and  $t_{m+1}$ . Note that since  $t_m$  and  $t_m$  can vary independently, the flux anomaly  $F_1$  modeled in terms of the cycle decay rate has twice the temporal degrees of freedom contained in the  $F_2$  anomaly series, for which adjacent cycle lengths both depend on the epoch  $t_m$  of the intervening minimum.

Applying an arbitrary smoothing to the solar record can considerably alter its climatic impact [Kelly and Wigley, 1992]. We choose instead to model the effect of convective heat-flux variability on TSI in terms of a first-order autoregressive (AR-1) process:

$$dW/dt + W/\tau = k F(t); (3)$$

here W is the solar irradiance anomaly, F(t) is the anomalous convective heat transport derived from either the cycle decay-rate (CD) model (Eq. 1) or the cycle length (CL) model (Eq. 2), k is an unknown proportionality constant, and  $\tau$  is a relaxation time for the convective anomaly. The relaxation time  $\tau \sim L^2/\nu$  was estimated by choosing the kinematic eddy viscosity  $\nu$  of the convective zone near the lower limit of a plausible range [ $\nu \sim 10^{12} - 10^{13}$  cm<sup>2</sup>sec<sup>-1</sup>, cf. Zeldovich et al. 1983] and the relevant length scale L as the depth of the convective zone (about 1/3 of the solar radius); this yields  $\tau = 12.6$  yr.

TSI anomalies were calculated from Lassen and Friis-Christensen's [1995] epochs for  $t_{\rm m}$  and  $t_{\rm M}$  (their Tables 1 and 2), using both the CD and CL formulations for the convective heat flux anomaly. Due to the linearity of the AR-1 process, the proportionality constants  $k_1$  and  $k_2$  in Eqs. (1) and (2) can be combined with the constant  $k_1$  in Eq. (3) into a single unknown scaling factor for

each of the irradiance curves, which are plotted in Fig. 1 in arbitrary units. The cycle-length (CL) model irradiance has a clear upward trend over the 140-year span shown in the figure, while the cycle decay (CD) rate model shows a sharp TSI increase between 1880 and 1937, with a subsequent decrease to values below the 140-year average. The implications of these characteristics of the modeled TSI for Sun-climate relations are explored next.

## Global Mean Temperature Response

We computed the response of the global mean temperature  $T_S$  to changes in radiative forcing by applying a simplified version of the IPCC upwelling-diffusion model [Kattenberg et al., 1996]; no distinction is made in this model version between land and ocean or northern and southern hemispheres. A fraction  $\Pi = 0.2$  of the temperature change is assumed to be downwelled by the thermohaline circulation, which is characterized by a fixed, globally averaged upwelling velocity of 4 m yr<sup>-1</sup>. Ensembles of runs were performed starting from an assumed zero temperature anomaly in 1850, in which the climate sensitivity S (defined as the equilibrium increase in  $T_S$  for doubled CO<sub>2</sub>) and solar forcing amplitude  $I_O$  (defined as the difference between the minimum and maximum values given by the irradiance curves in Fig. 1) were varied over fixed ranges. The model results for  $T_S$  were evaluated in terms of the percentage of variance accounted for in the Jones et al. [1994] reconstruction of  $T_S$  that spans the pre-Pinatubo instrumental temperature record (1854-1991). Possible effects of the "cold start" [e.g., Hasselmann et al., 1993] were addressed by examination of runs initialized in 1765 with IPCC-estimated anthropogenic forcing; the impact of this initialization on the model's temperature variation during the instrumental period was minimal.

We first performed a series of experiments using only the anomalous solar forcing given by the cycle-length (CL) TSI, with amplitudes ranging up to  $I_0 = 1.80$  Wm<sup>-2</sup>; at the same time, the model's climate sensitivity was varied from 0.5 to 5.0 °C (Fig. 2a, upper panel). The highest amount of variance in the *Jones et al.* [1994]  $T_s$  record that can be accounted for by the CL proxy model is 55%; it was obtained for a net irradiance amplitude of  $I_0 = 0.90$  Wm<sup>-2</sup>, which (assuming a planetary albedo of 30%) corresponds to a variation in the solar "constant" of 0.38% during the

last century. Lean et al. [1992] used linear regression of solar Ca II (HK) emission with respect to satellite-measured irradiance to estimate the TSI deficit for a noncycling state (such as the Maunder Minimum, ca. 1700) as 0.24%, which scales to a peak-to-peak variation of  $\sim$ 0.14% during the last century; Zhang et al.'s [1994] corresponding estimates, based on brightness changes in a sample of sunlike stars, span a range from about this value to an irradiance change over the last century of  $\sim$ 0.5%. Although consistent with a limited stellar sample, therefore, the amplitude of the CL-solar forcing required to fit the observed  $T_S$  record is considerably in excess of Dec-Cen variability inferred from the solar HK-irradiance relationship.

The best-fit climate sensitivity for the CL solar-only case (S = 5.0 °C) also exceeds plausible estimates for this parameter (1.5 - 4.5 °C, cf. *IPCC*, 1996; note that values of S within the IPCC range would imply greater TSI amplitude). The simulated  $T_S$  variation using the optimal ( $I_O$ , S) pair roughly matches the Jones et al. record until 1976 (Fig. 2a, lower panel), but even with these relatively large solar forcing and sensitivity parameters, the climate model is unable to capture the subsequent rapid increase which occurred prior to the Pinatubo eruption in 1991. This latter part of the record is expected *a priori* to be subject to the strongest anthropogenic effects; the failure of the simulated temperature to match the recent warming lends further support to the idea that this warming is not due to solar forcing alone.

To assess the relative contribution of human activities to the observed warming, we performed another series of model runs including the estimated radiative effects of greenhouse gases and sulfate emissions [IPCC, 1996], with the latter scaled to produce a global-mean forcing of -0.6 Wm<sup>-2</sup> in 1990. Note that the temporal profile of the net anthropogenic forcing (dotted line in Fig. 1) is similar to the solar irradiance profile generated by the CL model (with a correlation coefficient r = 0.57); both show an overall upward trend during the last century. The anthropogenic signature in the global mean temperature record, therefore, will resemble that of the CL-derived irradiance, making it difficult to distinguish their effects on the observed Dec-Cen variations of  $T_{S}$ .

The addition of anthropogenic forcing to the CL-derived solar irradiance increases the maximum  $T_s$  variance accounted for to 72% (Fig. 2b, upper panel), with the recent warming, in

particular, now fully captured (Fig. 2b, lower panel). The best-fit climate sensitivity has been reduced to 2.1  $^{\circ}$ C, well within the IPCC range. The implied solar forcing of 0.65 W/m² (or TSI variation of 0.27% during the last century) is also well within *Zhang et al.'s* [1994] estimated range for stellar variability, although it is still about twice the amplitude inferred from *Lean et al.'s* [1992] solar irradiance-HK relation. The "trade-off" between the effects of the CL-solar and anthropogenic forcing manifests itself as an extended, hyperbolic-shaped region in the  $I_0$ -S plane, for which the explained  $T_S$  variance is nearly the same as that obtained with the optimal parameters. The red-shaded area in Fig. 2b (upper panel), in particular, shows that the observed  $T_S$  record is consistent with relatively large CL-solar variations ( $I_0 > 1.2 \text{ W/m}^2$ , or ~0.5% TSI) during the last century, provided the climate sensitivity is restricted to values below the IPCC range ( $S < 1.5 ^{\circ}$ C). While these results indicate that the  $T_S$  record cannot be explained by CL-derived irradiance alone, therefore, they do not rule out a large role for solar forcing in producing the warming observed over the last century.

To examine the dependence of solar-forcing effects on the proxy model used to infer Dec-Cen TSI variability, we repeated the above calculations using the cycle decay-rate (CD) irradiance profile (solid line in Fig. 1). For the solar-forcing-only experiments the distribution of associated  $T_S$  variance in the  $I_0$ -S plane (Fig. 3a, upper panel) is quite similar to that obtained with the CL model, although its magnitude is diminished by about one-half. As for the CL model, the optimal temperature history simulated with CD-solar forcing reproduces the observed  $T_S$  variation fairly well until about 1976 (Fig. 3a, lower panel). Whereas the CL-derived irradiance shows a substantial increase since about 1970, however, the CD profile decreases monotonically from 1937 to 1980, with only a modest subsequent recovery (Fig. 1). The climate model responds to the CD-solar forcing with a steady temperature decrease since mid-century, causing large deviations from the observed  $T_S$  record during the recent warming.

The addition of anthropogenic forcing to the CD irradiance more than doubles the associated  $T_s$  variance (Fig. 3b, upper panel), reaching the same maximum (72%) obtained with the combined anthropogenic and CL forcing. The bulk of the improvement evidently comes from the last part of

the record, where the cooling produced by the CD solar-only forcing (Fig. 3a, lower panel) has been replaced by an accelerated warming (Fig. 3b, lower panel) which fits the observed upward trend of  $T_S$  within observational error [cf. *Jones et al.*, 1997]. The best-fit solar forcing has been reduced from 0.67 Wm<sup>-2</sup> for the combined CL case to 0.52 Wm<sup>-2</sup>, reflecting the poorer match between declining CD irradiance and the recent warming. The corresponding Dec-Cen TSI variability (0.22%) is roughly consistent with *Lean et al.'s* [1992] estimate of the maximum possible Maunder irradiance deficit.

#### **Discussion and Conclusions**

The relative roles of solar and anthropogenic forcing in producing climate change over the instrumental era, as deduced from the CL and CD experiments, are summarized in Table 1. For both irradiance proxies, the model results indicate that the steep rise in temperature observed in the early part of the century (1910-1940; see *Ghil and Vautard* [1991]) was largely caused by Dec-Cen solar variability. The warming documented in the full (1854-1991) instrumental  $T_S$  record is dominated by anthropogenic effects, however, with the CL-derived irradiance accounting for only 21% of the over-all temperature increase, while the CD-solar forcing actually offsets the warming over this time interval by a small amount (3%). Since the same anthropogenic forcing (dotted line in Fig. 1) was prescribed for all experiments, the climate sensitivity S required to fit the observed temperature record is greater, by almost 50%, for the CD irradiance profile.

The choice of a proxy irradiance model, therefore, can strongly influence the inferred role of solar forcing in recent climate change, as well as the quasi-equilibrium sensitivity of the terrestrial climate system deduced from the instrumental temperature record [see also Wigley et al., 1997]. Due to the near "orthogonality" (r = 0.07) of the CD-solar and anthropogenic forcing profiles (cf. Fig. 1), in particular, the subset of ( $I_0$ , S) values which is compatible with the observed  $T_S$  record forms a more compact region in parameter space than was obtained for the CL irradiance (compare upper panels of Figs. 2b and 3b). The combination of very high values of solar forcing ( $I_0 > 1.2$ 

Wm<sup>-2</sup>) and very low values of climate sensitivity (S < 1.5 °C) found to be compatible with the CL profile is thus effectively excluded by the CD irradiance model.

To the extent that the CD model more directly incorporates the physical processes responsible for Dec-Cen irradiance variability, therefore, these findings underscore the inability of solar forcing alone to explain recent temperature increases [see also *Solanki and Fligge*, 1998]. Since both irradiance models (combined with anthropogenic forcing) account for nearly the same amount of year-to-year  $T_S$  variance (72%), however, the results presented here cannot be used to ascertain which proxy TSI series best captures Dec-Cen solar variability. An accurate reconstruction of solar effects on global climate requires the development of robust and reliable proxy models which are related as directly as possible to the underlying mechanisms. The identification of turbulent convective processes which affect heat transport into the photosphere, as well as solar-cycle and spot decay rates, seems a promising approach [*Hoyt and Schatten*, 1993, 1997]. Further refinements in both its theoretical and observational aspects will be needed, however, to accurately constrain the solar contribution to recent climate change.

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Solar Model	Century-Scale Irradiance Change	Implied Climate Sensitivity	Solar Contribut 1854–1991	ion to Warming 1910–1940
CL	0.27%	2.1°C	21%	70%
CD	0.22%	3.0°C	-3%	58%

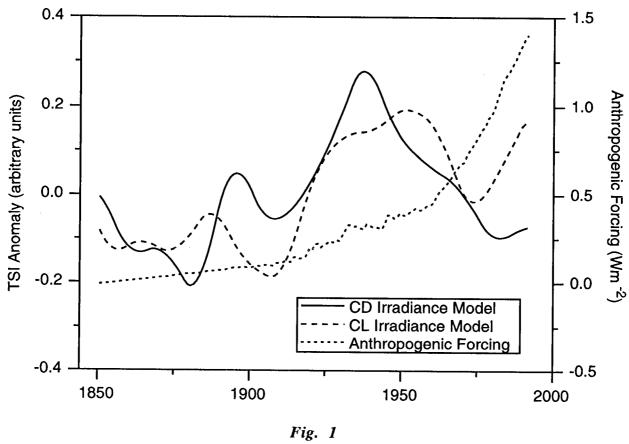
Table 1. Results from experiments using combined anthropogenic and solar forcing.

#### Figure Captions

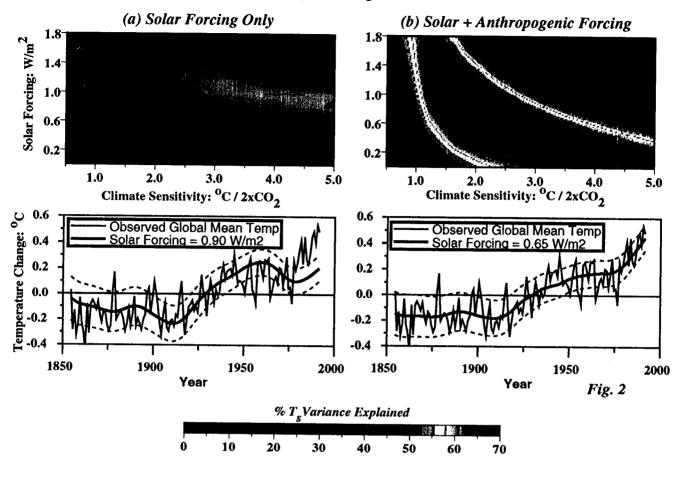
Fig. 1. Time series of radiative-forcing terms used to drive the climate model. TSI anomalies derived from solar cycle decay (CD) rate and solar cycle length (CL) proxies are shown in arbitrary units (left ordinate), and the anthropogenic forcing is shown in Wm<sup>-2</sup> (right ordinate).

Fig. 2. (a) Results of climate simulations using the anomalous TSI profile given by the solar cycle length (CL) proxy. The upper panel shows the percentage of variance in the Jones et al. (1994) yearly series of global mean surface temperature associated with the simulated  $T_S$  series, using the climate sensitivity ( $I_0$ ) and solar scaling (S) parameters defined by the abscissa and ordinate, respectively. The lower panel shows the corresponding temperature history obtained using the ( $I_0$ , S) pair that yields the highest variance; note that since the climate model simulates only temperature changes, both the simulated and observed series (solid red and blue curves, respectively) have been adjusted to have zero mean over the time interval plotted. The dashed red lines show  $2\sigma$  confidence intervals derived by *Jones et al.* [1997] for decadal temperature errors, and are plotted to indicate a plausible range of observational deviations from a "perfect" model. (b) as in (a), with the anthropogenic radiative forcing added to the CL-solar forcing.

Fig. 3. (a) As in Fig. 2a, using the anomalous TSI profile given by the solar cycle decay (CD) rate proxy. (b) As in Fig. 3a, with the anthropogenic radiative forcing added to CD-solar forcing.



# Solar Cycle Length Model



# **Solar Cycle Decay Model**

